

Extinction in metapopulations with environmental stochasticity driven by catastrophes

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- $\mathbf{X} = [X \ Y]^T$. The state of the metapopulation:
 - X is the number of suitable patches;
 - Y is the number of occupied patches.
- r, c, e, γ, p . Rates and other parameters.

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- Each occupied patch produces migrants at rate c , which may colonise empty, suitable patches

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- Each local population goes extinct at rate e

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on $S = \{(x, y) \mid x, y \in \mathbb{N}, 0 \leq y \leq x \leq N\}$.

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$$(x, y) \rightarrow (x - (i + j), y - j) \quad \text{at rate}$$

$$\gamma \binom{x - y}{i} \binom{y}{j} p^{i+j} (1 - p)^{x-i-j}.$$

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Finite State-space Processes

When N is finite, we can hope to evaluate measures of interest directly.

- Extinction times (a.k.a. first passage or 'hitting' times) are almost surely finite!
- Quasi-stationary distributions exist!
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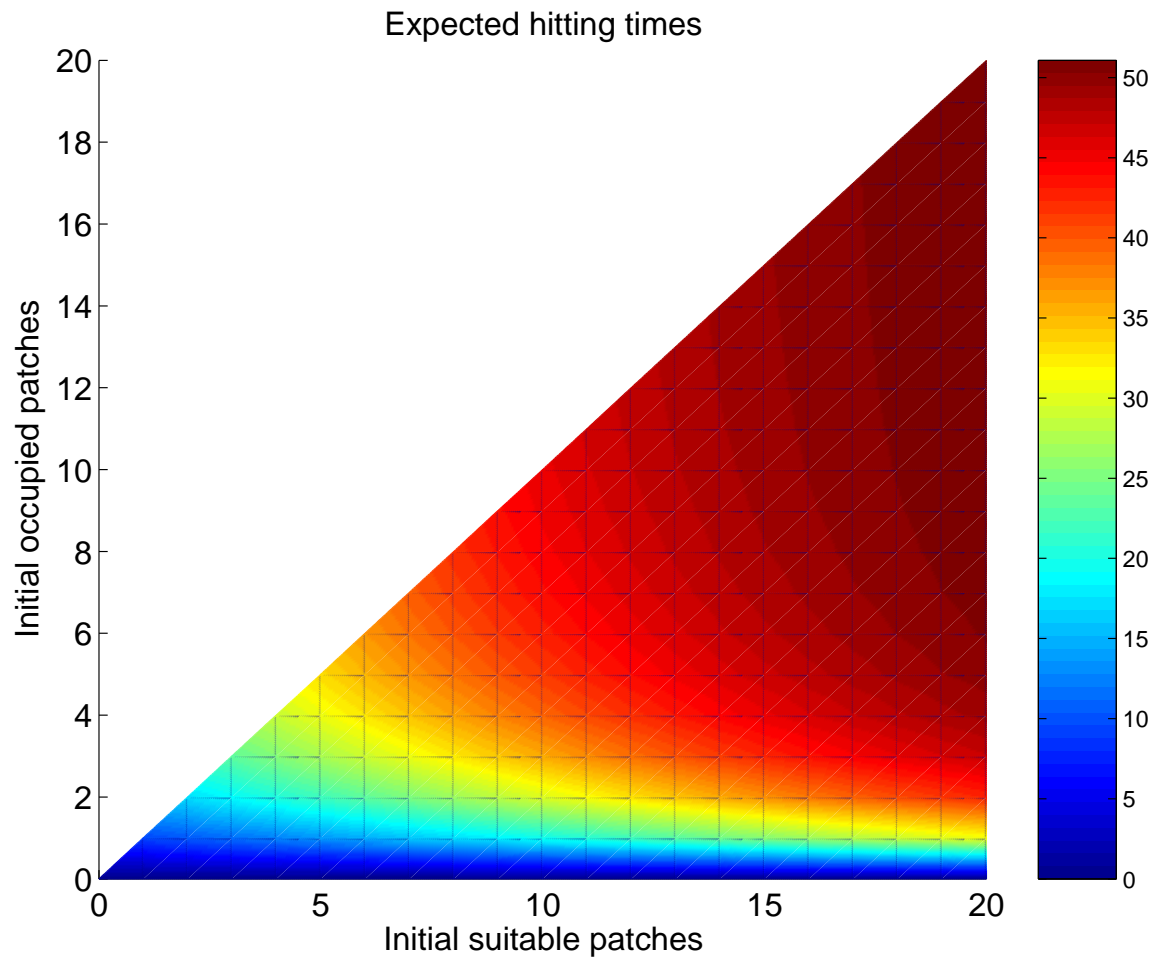
If these are easy to calculate (or approximate), we can use them to analyse the system.

Extinction Times

$$Q_C \tau = -1$$

(Generally, take the minimal,
non-negative solution.)

Extinction Times

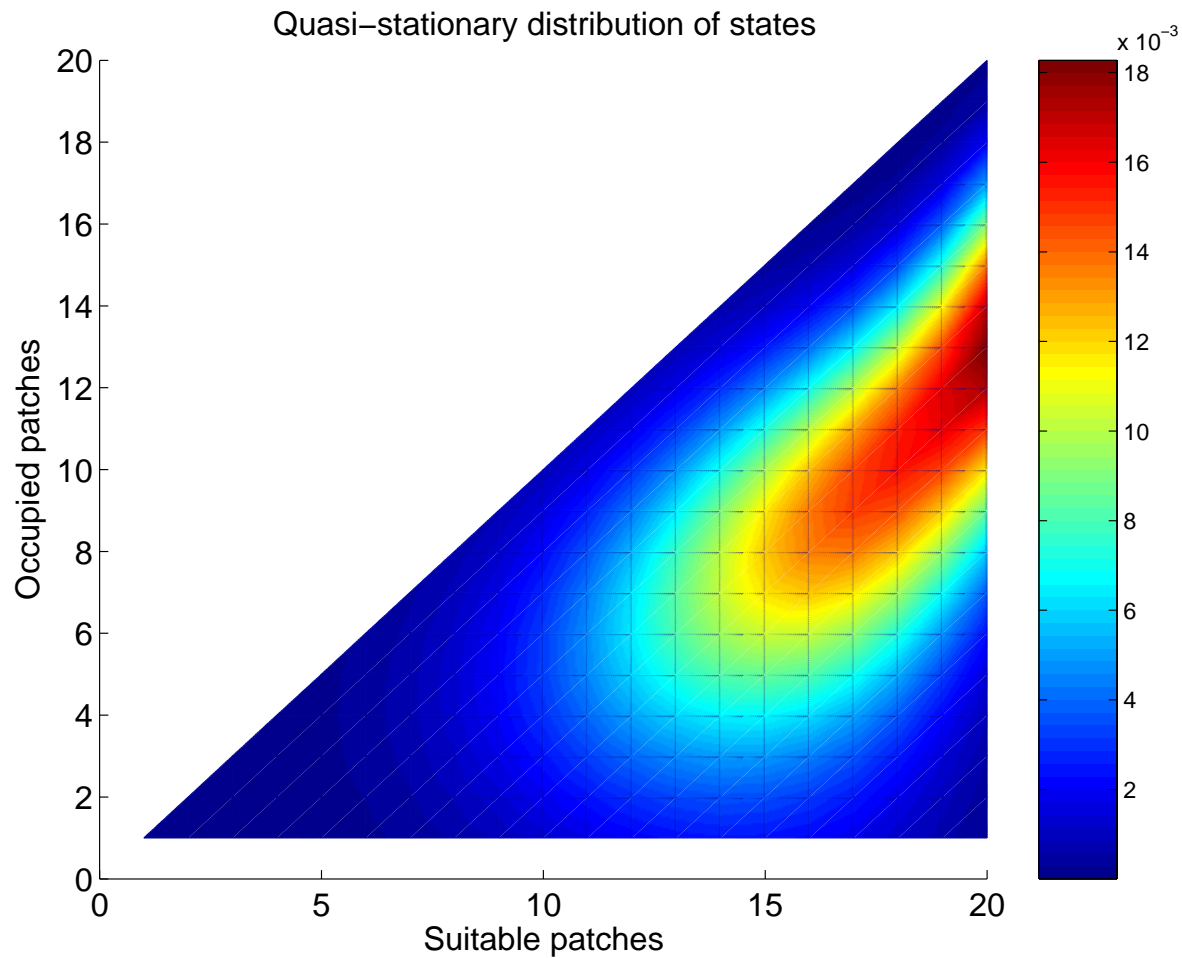


Quasi-stationary distributions

$$Q_C \mathbf{m} = -\lambda \mathbf{m}$$

(λ is the eigenvalue with max. real part.)

Quasi-stationary distributions



Difficulties as $N \uparrow$

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- To make progress, we need good approximations: e.g. stochastic differential equations for the limit as $N \rightarrow \infty$?

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$$\mathbf{a}(\mathbf{U}) = \begin{bmatrix} \partial u / \partial t \\ \partial v / \partial t \end{bmatrix} = \begin{bmatrix} r(1 - u) \\ cv(u - v) - ev \end{bmatrix},$$

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Catastrophes in the Limit

Treating catastrophes as a separate component...

- The arrival rate of catastrophes is unaffected by scaling.
- As $N \rightarrow \infty$, if T_1 is a catastrophe time,

$$\frac{\mathbf{X}(s, T_1)}{N} \xrightarrow{P} (1 - p)\mathbf{U}(s, T_1-).$$

A Stochastic Integral Equation

The limiting, scaled process:

$$d\mathbf{U}(s, t) = \mathbf{a}(\mathbf{U}(s, t))dt + \int_{\mathbf{M}} \mathbf{c}(\mathbf{U}(s, t), \mathbf{m}) \mathcal{P}[d\mathbf{m}, dt; \gamma]$$

- $\mathbf{c}(\mathbf{U}, d\mathbf{m})$ describes effect of catastrophes
- Poisson random measure \mathcal{P} describes arrival of catastrophes and their magnitudes, \mathbf{m} .
- Generalised Itô formula gives first passage times. (Gihman & Skorohod, 1972)

First Passage Times

First passage times, $\tau_G(\mathbf{U}_0)$, into a closed set $S \setminus G$ (i.e. out of G), starting from \mathbf{U}_0 , are a twice continuously differentiable solution, $g(\mathbf{U})$, to

$$(Lg)(\mathbf{U}) = -1, \mathbf{U} \in G$$

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- In the present case, $(Lh)(\mathbf{U})$ is given by

$$(Lh)(\mathbf{U}) = \nabla h(\mathbf{U}) \cdot \mathbf{a}(\mathbf{U}) - \gamma h(\mathbf{U}) + \gamma h((1-p)\mathbf{U}).$$

Solving First Passage Times

Slightly different conditions:

- $g(\mathbf{U})$ should be *continuous* along all trajectories $\mathbf{U}(s, t)$, and piecewise smooth along other smooth paths.
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Solve in 'steps': G_n is the region from which *at least* n catastrophes are needed to leave G .

Solving First Passage Times

The solution has the form

$$e^{-\gamma t} g(\mathbf{U}(s, t)) = - \int_0^t \gamma e^{-\gamma r} g((1-p)\mathbf{U}(s, r)) dr \\ - \gamma^{-1} [1 - e^{-\gamma t}] + C_1(s),$$

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$$C_1(s) = \int_0^\infty \gamma e^{-\gamma r} g((1-p)\mathbf{U}(s, r)) dr + \gamma^{-1}.$$

Solving First Passage Times

Hence (along trajectories that remain within G)

$$g(\mathbf{U}(s, t)) = \frac{1}{\gamma} + e^{\gamma t} \int_t^{\infty} \gamma e^{-\gamma r} g((1-p)\mathbf{U}(s, r)) dr.$$

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- if $G = G_1$, $g(\mathbf{U}) = \gamma^{-1}$ for all \mathbf{U} on trajectories remaining in G ;
- if $\mathbf{U}_{\infty} = \lim_{t \rightarrow \infty} \mathbf{U}(s, t)$ is in G , then $g(s, \infty-) = \gamma^{-1} + g((1-p)\mathbf{U}_{\infty})$.

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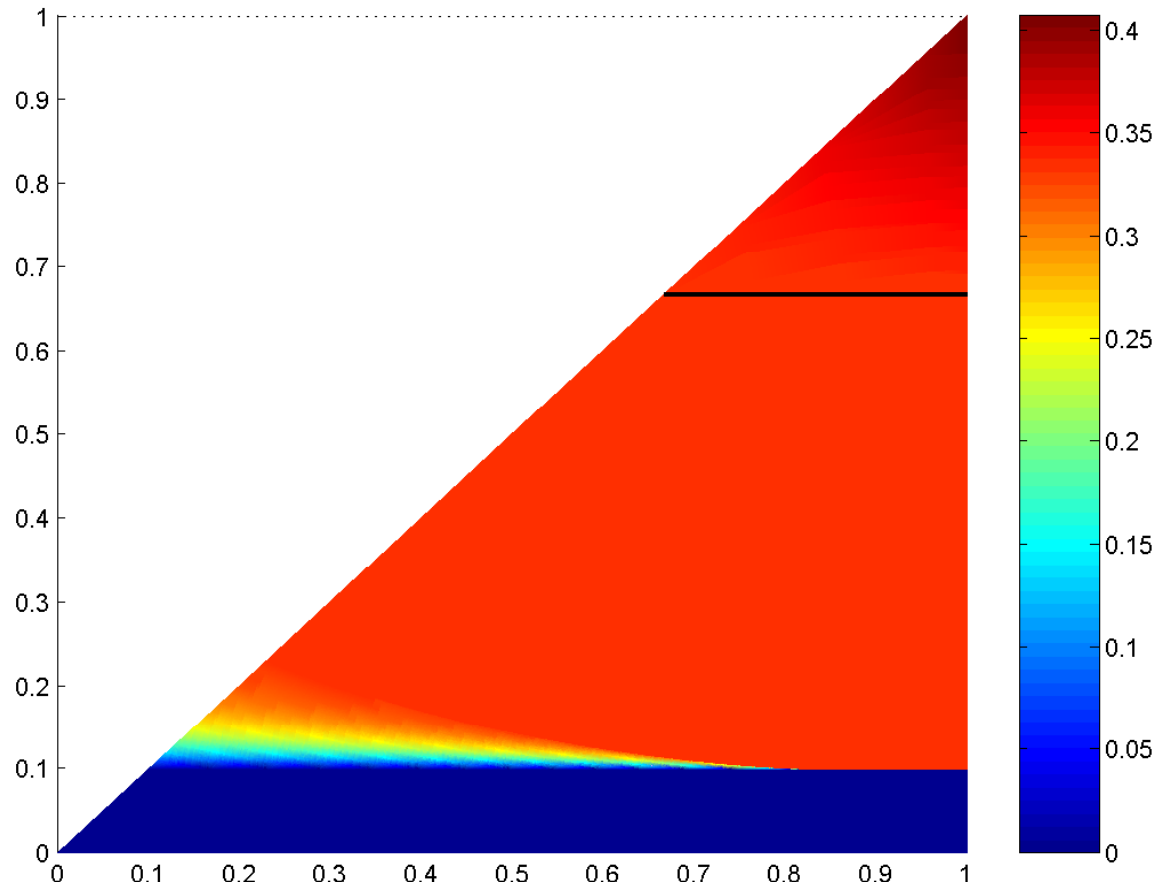
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- $g(\mathbf{U}(s, t)) = \gamma^{-1}$, for all trajectories not leaving G_1 in finite time,
- solutions for trajectories heading out of G using the deterministic hitting time and a truncated exponential law, and
- a system of DEs $[\partial u/\partial t, \partial v/\partial t, \partial g/\partial t]$ gives first passage times for trajectories starting on higher steps.

Solutions: A Special Case



Solutions: General Case

The general case is a little more difficult.

- Recall that

$$g(\mathbf{U}(s, t)) = \frac{1}{\gamma} + e^{\gamma t} \int_t^{\infty} \gamma e^{-\gamma r} g((1-p)\mathbf{U}(s, r)) dr.$$

Solutions: General Case

The general case is a little more difficult.

- Define a mapping $K : H \rightarrow H$,

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with H being the set of bounded functions $f : G \rightarrow \mathbb{R}_+$ under the condition

$$f(\mathbf{U}(s, t)) \geq \frac{1}{\gamma} + e^{\gamma t} \int_t^{\infty} \gamma e^{-\gamma r} f((1-p)\mathbf{U}(s, r)) dr.$$

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- $f \geq K(f)$, $f \in H$, so we might hope that the iterative application of K would lead to a fixed point, but...
- $H \neq \emptyset$ is equivalent to the existence of a solution, h , to

$$(Lh)(\mathbf{U}) \leq -1, \mathbf{U} \in G$$

$$h(\mathbf{U}) \geq 0, \mathbf{U} \notin G,$$

\equiv to a condition from Gihman & Skorohod for the existence of a solution $\tau_G \leq h$.

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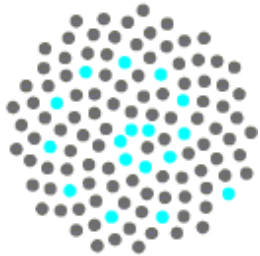
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 - (i) take $G' \supset G$ so that the first passage out of G' only depends on u ;
 - (ii) find $h(u) = \tau_{G'}(u)$ (using H & T);
 - (iii) then $h(u)$ satisfies the inequality condition for all v such that $(u, v) \in S$.

Thanks

- Phil Pollett and Hugh Possingham (advisors),
Chris Wilcox and Josh Ross.



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