

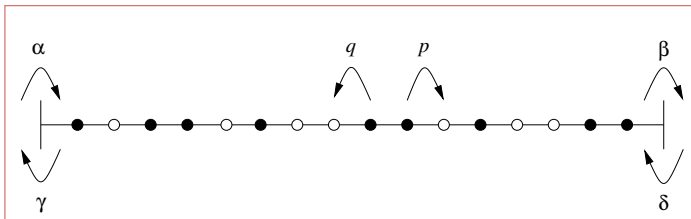
The asymmetric exclusion process and Askey-Wilson polynomials

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Workshop on Stochastics and Special Functions
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Partially Asymmetric Simple Exclusion Process (PASEP)

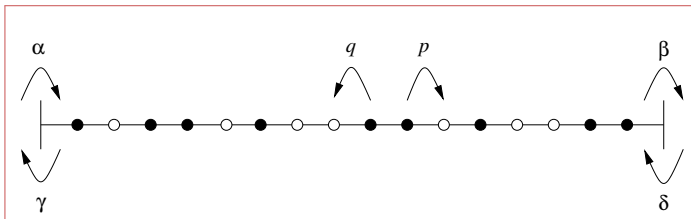


- Probability to find a configuration \mathcal{C} at time t : $P(\mathcal{C}, t)$
- Time evolution given by the Master Equation:

$$\frac{\partial}{\partial t} P(\mathcal{C}, t) = \sum_{\mathcal{C}' \neq \mathcal{C}} P(\mathcal{C}', t) W(\mathcal{C}' \rightarrow \mathcal{C}) - P(\mathcal{C}, t) \sum_{\mathcal{C}' \neq \mathcal{C}} W(\mathcal{C} \rightarrow \mathcal{C}')$$

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Denote $\mathcal{C} = (\tau_1, \dots, \tau_N)$, $\tau_i \in \{0, 1\}$ and use unnormalised weights f :

$$P(\mathcal{C}, t) = \frac{1}{Z_N} f(\tau_1, \dots, \tau_N)$$

Theorem

\exists (infinite) matrices D, E , a row vector ℓ and column vector r such that for the stationary state ($t = \infty$):

$$f(\tau_1, \dots, \tau_N) = \ell \cdot \prod_{i=1}^N [\tau_i D + (1 - \tau_i) E] \cdot r$$

Example:

$$f(10110100) = \ell \cdot \text{DEDDEDEE} \cdot r$$

Normalisation:

$$Z_N = \ell \cdot (D + E)^N \cdot r = \ell \cdot C^N \cdot r, \quad C = D + E$$

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If matrices are known, it is easy to compute quantities of interest.

■ Density:

$$\rho_i = \mathbb{E}[\tau_i] = \frac{1}{Z_N} \ell \cdot C^{i-1} D C^{N-i} \cdot r$$

■ Current:

$$J_{i,i+1} = \mathbb{E}[\tau_i(1 - \tau_{i+1})] = \frac{1}{Z_N} \ell \cdot C^{i-1} D E C^{N-i-1} \cdot r$$

In the case of TASEP ($q = 0 = \gamma = \delta$), matrices satisfy the relations:

$$D E = D + E (= C),$$

$$D \cdot r = \frac{1}{\beta} r,$$

$$\ell \cdot E = \frac{1}{\alpha} \ell.$$

Hence

$$J_{i,i+1} = J = \frac{1}{Z_N} \ell \cdot C^{i-1} (D + E) C^{N-i-1} \cdot r = \frac{Z_{N-1}}{Z_N}$$

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Matrix algebra

Master equation can be written as

$$\frac{\partial f}{\partial t} = H f$$

with

$$H = h_0 + \sum_{i=1}^{N-1} h_{i,i+1} + h_N$$

Example:

$$h_{i,i+1} f(\dots 10 \dots) = q f(\dots 01 \dots) - f(\dots 10 \dots).$$

or

$$h_{i,i+1}(DE) = qED - DE.$$

Example:

$$h_0(D) = \alpha E - \gamma D.$$

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Proof of Theorem.

Stationary state:

$$f(\tau_1, \dots, \tau_N) = \ell \cdot \prod_{i=1}^N [\tau_i D + (1 - \tau_i) E] \cdot r$$

Let D and E satisfy the relations

$$\begin{aligned} -h_{i,i+1}(DE) &= DE - qED &= D + E \\ h_N(E) \cdot r &= (\beta D - \delta E) \cdot r &= r \\ \ell \cdot h_0(D) &= \ell \cdot (\alpha E - \gamma D) &= \ell. \end{aligned}$$

Then $Hf = (h_0 + \sum_{i=1}^{N-1} h_{i,i+1} + h_N) f$ telescopes to 0. □

Note: Different model \rightarrow same matrix product formalism but with different matrix algebra.

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Explicit representation for TASEP

In the case $q = \gamma = \delta$:

$$D = \begin{pmatrix} 1 & 1 & 0 & 0 & \cdots \\ 0 & 1 & 1 & 0 & \\ 0 & 0 & 1 & 1 & \\ 0 & 0 & 0 & 1 & \\ \vdots & & & & \ddots \end{pmatrix}, \quad E = \begin{pmatrix} 1 & 0 & 0 & 0 & \cdots \\ 1 & 1 & 0 & 0 & \\ 0 & 1 & 1 & 0 & \\ 0 & 0 & 1 & 0 & \\ \vdots & & & & \ddots \end{pmatrix}$$

$$l = \kappa(1, a, a^2, a^3, \dots), \quad r = \kappa(1, b, b^2, b^3, \dots)^T$$

where

$$a = \frac{1-\alpha}{\alpha}, \quad b = \frac{1-\beta}{\beta}, \quad \kappa = \frac{1}{\alpha} + \frac{1}{\beta} - \frac{1}{\alpha\beta}.$$

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$$C = D + E = \begin{pmatrix} 2 & 1 & 0 & 0 & \cdots \\ 1 & 2 & 1 & 0 & \\ 0 & 1 & 2 & 1 & \\ 0 & 0 & 1 & 2 & \\ \vdots & & & & \ddots \end{pmatrix}$$

Want to compute $Z_N = \ell \cdot C^N \cdot r \rightarrow$ diagonalisation:

$$C \cdot e(\theta)^T = 2(1 + \cos \theta) e(\theta)^T$$

where

$$e(\theta)^T = \frac{1}{\sin \theta} \begin{pmatrix} \sin \theta \\ \sin 2\theta \\ \sin 3\theta \\ \vdots \end{pmatrix}$$

Elements of $e(\theta)$ are Chebyshev polynomials of the second kind.

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Due to orthonormality of Chebyshev polynomials:

$$\mathbb{I} = \frac{1}{\pi} \int_0^{2\pi} \sin^2 \theta \mathbf{e}(\theta)^T \mathbf{e}(\theta) d\theta$$

Hence

$$\begin{aligned} Z_N &= \frac{1}{\pi} \int_0^{2\pi} \sin^2 \theta \ell \cdot C^N [\mathbf{e}(\theta)^T \mathbf{e}(\theta)] \cdot r \\ &= \frac{1}{\pi} \int_0^{2\pi} \sin^2 \theta \left[\ell \cdot \mathbf{e}(\theta)^T \right] 2^N (1 + \cos \theta)^N \left[\mathbf{e}(\theta) \cdot r \right] d\theta \end{aligned}$$

$$\ell \cdot \mathbf{e}(\theta)^T = \frac{\kappa}{\sin \theta} \sum_{n=0}^{\infty} a^n \sin(n+1)\theta = \frac{\kappa}{(1 - ae^{i\theta})(1 - ae^{-i\theta})}$$

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- Similar computations for **density**, **current** and other observables
- Asymptotic analysis \rightarrow non-analytic dependence on a and b
- Phase transitions

Next: allow backhopping: $q \neq 0$.

$$Z_N = \frac{1}{\pi} \int_0^{2\pi} \sin^2 \theta \frac{\kappa^2 2^N (1 + \cos \theta)^N}{(1 - a e^{i\theta})(1 - a e^{-i\theta})(1 - b e^{i\theta})(1 - b e^{-i\theta})}.$$

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$$DE - qED = D + E$$

$$\beta D \cdot r = r$$

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Diagonalise $C = D + E \rightarrow$ Find $e(\theta)$ such that

$$C \cdot e(\theta)^T = \frac{2(1 + \cos \theta)}{1 - q} e(\theta)^T, \quad e(\theta) = (e_n(\theta))_{n=0}^{\infty}$$

Three-term recurrence:

$$2\cos \theta \tilde{e}_n(\theta) = (1 - q^n)(1 - abq^{n-1})\tilde{e}_{n-1}(\theta) + (a + b)q^n \tilde{e}_n(\theta) + \tilde{e}_{n+1}(\theta)$$

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Three term recurrence for **Al-Salam-Chihara** orthogonal polynomials.

Orthogonality:

$$\frac{1}{2\pi} \frac{(q, ab; q)_{\infty}}{(q, ab; q)_n} \int_0^{\pi} \frac{(e^{2i\theta}, e^{-2i\theta}; q)_{\infty}}{(ae^{i\theta}, ae^{-i\theta}, be^{i\theta}, be^{-i\theta}; q)_{\infty}} \tilde{e}_n(\theta) \tilde{e}_m(\theta) d\theta = \delta_{n,m}$$

Normalisation:

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Same procedure results in three-term recurrence for **Askey-Wilson** orthogonal polynomials.

$$\tilde{e}_n(\theta) = \frac{(ab, ac, ad; q)_n}{a^n} \sum_{k=0}^{\infty} \frac{(q^{-n}, q^{n-1}abcd, ae^{i\theta}, ae^{-i\theta}; q)_k}{(ab, ac, ad, q; q)_k} q^k.$$

Normalisation:

$$Z_N = \frac{(q, ab, ac, ad, bc, bd, cd; q)_{\infty}}{2\pi(abcd; q)_{\infty}} \times \int_0^{\pi} \frac{(e^{2i\theta}, e^{-2i\theta}; q)_{\infty}}{(ae^{i\theta}, ae^{-i\theta}, be^{i\theta}, be^{-i\theta}, ce^{i\theta}, ce^{-i\theta}, de^{i\theta}, de^{-i\theta}; q)_{\infty}} \left[\frac{2(1 + \cos \theta)}{1 - q} \right]^N.$$

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Decay rates

Decay $P(\mathcal{C}, t) - P(\mathcal{C}, \infty) \propto e^{-\Lambda t}$.

$$\Lambda(\{z_j\}) = \alpha + \beta + \gamma + \delta + \sum_{j=1}^{L-1} \frac{(q-1)^2 z_j}{(1-z_j)(qz_j-1)}$$

$$\left[\frac{qz_j - 1}{1 - z_j} \right]^{2L} K(z_j) = \prod_{l \neq j}^{L-1} \left[\frac{qz_j - z_l}{z_j - qz_l} \right] \left[\frac{q^2 z_j z_l - 1}{z_j z_l - 1} \right]$$

with $K(z) = \tilde{K}(z, \alpha, \gamma) \tilde{K}(z, \beta, \delta)$ and

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Integral equation

Taking log gives:

$$Y_L(z) := g(z) + \frac{1}{L}k(a, b, c, d; z) + \frac{1}{L} \sum_{l=1}^{L-1} \log S(z_l, z)$$

Curve: $Y_L(z_j) = -\pi + \frac{2\pi j}{L}$.

Now use Cauchy:

$$\frac{1}{L} \sum_{j=1}^{L-1} h(z_j) = \oint_{C_1+C_2} \frac{dz}{2\pi i} h(z) \cot\left(\frac{1}{2}LY_L(z)\right)$$

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Given Y_L we can compute Λ_1 (and other quantities of interest). To find Y_L we expand

$$Y_L(z) = \sum_{n=0}^{\infty} L^{-n} y_n(z), \quad \xi = z_c + \sum_{n=1}^{\infty} L^{-n} (\delta_n + i \eta_n),$$

Full solution for $y_1(z)$:

$$\begin{aligned}
 y_1(z) = & -i \ln \left[-\frac{z}{z_c} \frac{1-z_c^2}{1-z^2} \right] + \kappa_1 - i \ln(ab) - \lambda_1 \ln(-z_c) \\
 & - i \ln \left[\frac{(-c/z; q)_\infty (-cz; q)_\infty (-z/a; q)_\infty (-qaz_c; q)_\infty}{(-c/z_c; q)_\infty (-cz_c; q)_\infty (-z_c/a; q)_\infty (-qaz; q)_\infty} \right] \\
 & - i \ln \left[\frac{(-d/z; q)_\infty (-dz; q)_\infty (-z/b; q)_\infty (-qbz_c; q)_\infty}{(-d/z_c; q)_\infty (-dz_c; q)_\infty (-z_c/b; q)_\infty (-qbz; q)_\infty} \right] \\
 & + \lambda_1 \ln \left[\frac{(qz_c/z; q)_\infty (qzz_c; q)_\infty^2}{(qz/z_c; q)_\infty (qz_c^2; q)_\infty^2} \right] + \lambda_1 \ln \left(\frac{z - z_c^{-1}}{z_c - z_c^{-1}} \right), \\
 z_c = & -1/\sqrt{ab}
 \end{aligned}$$

$$\Lambda_1 = -D_+ - D_- + 2\sqrt{D_+D_-} - \sqrt{D_+D_-}\frac{\pi^2}{L^2} + \mathcal{O}(L^{-3}).$$

with

$$D_{\pm} = (1 - q) \frac{\rho^{\pm}(1 - \rho^{\pm})}{\rho^+ - \rho^-}$$

given in terms of low and high density phase densities

$$\rho^- = \frac{1}{1+a}, \quad \rho^+ = \frac{b}{1+b}$$

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